# **STICM**

### Select / Special Topics in Classical Mechanics

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STiCM Lecture 15:

Unit 5 : Non-Inertial Frame of Reference *Real effects of pseudo-forces!*

#### **Unit 5: Inertial and non-inertial reference frames.**

**Moving coordinate systems. Pseudo forces. Inertial and non-inertial reference frames.** 

**Deterministic cause-effect relations in inertial frame, and their** *modifications* **in a non-inertial frame.**



**Six Flags over Georgia**

### **Learning goals:**

Understand "Newton's laws hold only in an inertial frame".

Distinction between an inertial and a non-inertial frame is linked to what we consider as a fundamental physical force/interaction.

Electromagnetic/electroweak, nuclear, gravity) / pseudo-force (centrifugal, Coriolis etc.) / Friction

We shall learn to interpret the 'real effects of pseudo-forces' in terms of Newtonian method. *Re-activate* 'causality' in noninertial frame of reference!

# Deterministic cause-effect relation in inertial frame,

# *and its adaptation in a non-inertial frame!*

The law of inertia enables us recognize an inertial frame of reference as one in which motion is *self-sustaining,* 

*determined entirely by initial* 

*conditions alone*.

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# **Just where is the inertial frame? Newton envisaged the inertial frame to be**

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**located in** 

*deep space, amidst* 

*distant stars***.**

Kabhi kabhi mere dil mein, Khayaal aata hein…..

Tu ab se pahile, sitaron mein bas rahi thi kahin,

Tujhe jamin pe, bulaya gaya hei mere liye….





Time t is the same in the red frame and in the blue frame.



First derivative with respect to time is different!

So what? Remember Galileo?

Second derivative with respect to time is, however, very much the same!



Essentially the same 'cause' explains the 'effect'

(acceleration) according to the same 'principle of causality'.

Laws of Mechanics : same ipcall INERTIAL FRAMES.

# $\overline{F}$  = ma = ma ' =  $F$ '

• A frame of reference moving with respect to an inertial frame of reference at constant velocity is also an inertial frame.

• The same force  $\left(\stackrel{\rightarrow}{F} = ma\right)$  explains the linear response (effect/acceleration is *linearly proportional to the cause/interaction)* relationship in all inertial frames.

Galileo's experiments that led him to the law of inertia.





**Galileo Galilee** 1564 - 1642



 $\vec{F} = m\vec{a}$  Linear Response. *Effect* is proportional to the *Cause* Principle of causality.

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 $\vec{F} = m\vec{a}$  Linear Response

$$
\vec{W} = m\vec{g}
$$

Weight = Mass  $\times$  acceleration due to gravity

# Lunatic exercise!









# Another lunatic exercise! *Is it easier to stop a charging bull on the Moon ?*

### **Ian Usmovets Stopping an Angry Bull (1849)**

Evgraf Semenovich Sorokin (born as Kostroma Province) 1821-1892. Russian artist and teacher, a master of historical, religious and genre paintings. PCD\_STiCM

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What do we mean by a CAUSE?

CAUSE is that physical agency

which generates an **EFFECT** !

### EFFECT: departure from equilibrium

CAUSE: 'forCE' 'intErACtion'

### fundamental interaction

### Electromagnetic, gravitational nuclear weak/strong

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Galilean relativity

 $\overrightarrow{F} = m\overrightarrow{a} = m\overrightarrow{a} = \overrightarrow{F}$ 

- A frame of reference moving with respect to an inertial frame of reference at constant velocity is also an inertial frame.
- The same force  $(F=ma)$  explains the linear response (effect/acceleration is linearly proportional to the cause/interaction) relationship in all inertial frames.
- "An inertial frame is one in which Newton's laws hold"

Galilean relativity Galilean relativity red frame and in the green, double-primed frame.

# Physics in an accelerated frame of reference.

# What happens to the (Cause, Effect) relationship?

What happens to the Principle of Causality / Determinism? PCD STICM







# **Same cause-effect relationship** does not explain the dynamics.





Laws of Mechanics are *not* the same in a FRAME OF REFERENCE that is accelerated with respect to an inertial  $\overrightarrow{OO}$   $\cdot$  =  $\overrightarrow{u}_i t + \frac{1}{2} \overrightarrow{f} t^2$ <br>Laws of Mechanics are<br>in a FRAME OF REFEF<br>accelerated with respectrame.

$$
\vec{F}'' = \vec{F} - \vec{F}_{pseudo}
$$

The force/interaction which explained the acceleration in the inertial frame does not account for the acceleration in the accelerated frame of reference.

$$
\vec{F}'' = \vec{F} - \vec{F}
$$
 *pseudo* Time t is the same in the red frame and in the green, double-primed frame.  
\n
$$
\vec{F}'' = \vec{F}
$$
 *real* / *physical* -  $\vec{F}$  *pseudo*

$$
\begin{array}{cccc}\n\cdot & & & \\
\cdot & & & \\
\end{array}
$$

Real effects of pseudo-forces!

*P. Chaitanya Das, G. Srinivasa Murthy, Gopal Pandurangan and P.C. Deshmukh* Resonance, Vol. 9, Number 6, 74-85 (2004) *http://www.ias.ac.in/resonance/June2004/pdf/June2004Classroom1.pdf*

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Deterministic / cause-effect / relation holds only in an inertial frame of reference.

This relation inspires in our minds an intuitive perception of a fundamental interaction / force (EM, Gravity, Nuclear strong/weak).

*Adaptation of causality in a non-inertial frame requires 'inventing' interactions that do not exist.*

*– for these 'fictitious forces', Newton's laws are of course not designed to work!* 

Galilean relativity:  $\vert \vert$  Time t is the same in all frames of references.

### RELATED ISSUES:

**Weightlessness** 

 What is Einstein's weight in an elevator accelerated upward/downward?



Sergei Bubka (Ukrainian: Сергій Бубка) (born December 4,

1963) is a retired Ukrainian pole vaulter. He represented the

Soviet Union before its dissolution in 1991. He is widely

regarded as the best pole vaulter ever.

*Reference: http://www.bookrags.com/Sergei\_Bubka* <u>PCD\_STiCM</u>















http://www.stabhoch.com/bildserien/20030825\_Isinbajeva\_465/bildreihe.html







**2008 Olympics champion: 5.05meters What enables the pole vaulting champion, Yelena Isinbayeva, twist her body in flight and clear great**  PCD**\_STGM**ts?

We will take a Break… *…… Any questions ?*

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**Next L16 : Non-Inertial Frame of Reference** 'cause', where there isn't one! Real effects of pseudo-forces!

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#### STiCM Lecture 16: Unit 5

Non-Inertial Frame of Reference 'cause', where there isn't one! Real effects of pseudo-forces!



#### Solid

**Holds Shape** 

**Fixed Volume** 



#### Solid

Liquid

of Container



#### http://www.lerc.nasa.gov/WWW/K-12/airplane/state.html





Glenn Research Center



 What will be the shape of a tiny little amount of a liquid in a closed (sealed) beaker when this '**liquid-in-a-beaker**' system is (a) on earth (b) orbiting in a satellite around the earth.

In a **liquid** the molecular forces are weaker than in a solid.

A liquid takes the shape of its container with a free surface in a gravitational field.

#### **Regardless of gravity, a liquid has a fixed volume.**

Other than GRAVITY, are there other interactions that could influence the shape of the liquid's free surface?



Intra/Inter molecular forces provide the concave or the convex meniscus to the liquid. These interactions become dramatically consequential in microgravity<sub>STiCM</sub>

The actual shape an amount of liquid will take in a closed container in microgravity depends on whether the adhesive, or the cohesive, forces are strong.

Accordingly, the liquid may form a floating ball inside, or stick to the inner walls of the container and leave a cavity inside!

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You can have zerogravity experience by booking your zero-gravity flight!





#### <http://www.gozerog.com/>

Flight ticket (one adult) : little over \$5k



### candle flame on Earth

On Earth, gravity-driven buoyant convection causes a candle flame to be teardrop-shaped and carries soot to the flame's tip, making it yellow.

http://www.grc.nasa.gov/WWW/RT/RT1996/6000/ 6726f.htm

http://science.nasa.gov/headlines/y2000/ast12ma y\_1.htm 29/09/09



34 In microgravity, where convective flows are absent, the flame is spherical, soot-free, and TICM<br>DIUE.



 $1$ <sup>11</sup><sup>2</sup> 2  $G\frac{m_1m_2}{2}$ *r*

# There is no 'insulation' from gravity!

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Galilean relativity

# Observations in a rotating frame of reference.


reference are different from those in the rotating frame of reference.

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We shall ignore translational motion of the new frame relative to the inertial frame  $F_{\rm I}$ .

We have already considered that.

We can always superpose the two (translational and rotating) motions.



Clearly, its time-derivative in the rotating frame is zero, but not so in the inertial frame.

$$
\frac{d}{dt} \frac{d}{dt} \frac{d}{dt}
$$

$$
\left(\frac{\mathrm{d}}{\mathrm{d}t}\right)_R \vec{b} \neq \left(\frac{\mathrm{d}}{\mathrm{d}t}\right)_I \vec{b}
$$

The particle seen at 'rest' in the rotating frame would appear to have moved to a new location as seen by an observer in  $F_1$ .



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# Often, one uses the term 'SPACE-FIXED FRAME OF REFERENCE' for  $F_I$ , and 'BODY-FIXED FRAME OF REFERENCE' for  $F_R$ .

We shall develop our analysis for an arbitrary vector  $\,b\,$  , the only condition being that it is itself not a timederivative in the rotating frame of some another vector.

*No* vector 
$$
\vec{q}
$$
 exists such that  $\left(\frac{d}{dt}\right)_R \vec{q} = \vec{b}$   
 $\left(\frac{d}{dt}\right)_R \vec{b} = \vec{0}$  in the rotating frame  $F_R$ .  
*Question*: What is  $\left(\frac{d}{dt}\right)_L \vec{b}$ ?

$$
d\vec{b} = \vec{b}(t+dt) - \vec{b}(t) = |d\vec{b}| \hat{u}
$$
  
\nwhere  $\hat{u} = \frac{\hat{n} \times \hat{b}}{|\hat{n} \times \hat{b}|} \cdot \xi = \angle(\hat{n}, \hat{b})$   
\n
$$
|d\vec{b}| = (b \sin \xi)(d\psi) \qquad \text{b} \sin \xi
$$
  
\n
$$
d\vec{b} = (\vec{b} \sin \xi)(d\psi) \qquad \text{b} \sin \xi
$$
  
\nThese two terms are equal and hence cancel.  
\n
$$
d\vec{b} = d\psi \hat{n} \times \vec{b}
$$
  
\n
$$
d\vec{b} = (\vec{\omega} dt) \times \vec{b}
$$
  
\n
$$
d\vec{b} = (\vec{\omega} dt) \times \vec{b}
$$
  
\n
$$
\text{since } \vec{\omega} = \frac{d\psi}{dt} \hat{n} \qquad \Rightarrow \left(\frac{d}{d\tau}\right) \vec{b} = \vec{\omega} \times \vec{b}
$$



Remember! The vector  $\vec{b}$  itself did not have any time-dependence in the rotating frame.

If  $\overline{b}$  has a time dependence in the <u>rotating frame</u>, the following operator equivalence would follow:

$$
\left(\frac{d}{dt}\right)_I \vec{b} = \vec{\omega} \times \vec{b} + \left(\frac{d}{dt}\right)_R \vec{b}
$$
\n
$$
\text{Operator} \quad \left(\frac{d}{dt}\right)_I = \left(\frac{d}{dt}\right)_R + \left(\frac{d}{\omega} \times \frac{d}{dt}\right)_L \text{ where } \vec{b} \text{ is the same value.}
$$

$$
\left(\frac{d}{dt}\right)_I = \left(\frac{d}{dt}\right)_R + \overrightarrow{\omega} \times
$$

$$
\left(\frac{d}{dt}\right)_I \vec{r} = \left(\frac{d}{dt}\right)_R \vec{r} + \vec{\omega} \times \vec{r}
$$

# Operating twice:

Operating twice:  
\n
$$
\left(\frac{d}{dt}\right)_l \left(\frac{d}{dt}\right)_l \vec{r} = \left(\frac{d}{dt}\right)_R \left\{ \left(\frac{d}{dt}\right)_k \vec{r} + \vec{\omega} \times \vec{r} \right\} + \left(\frac{d}{dt}\right)_l \vec{r} + \vec{\omega} \times \vec{r} \right\}
$$
\n
$$
+ \vec{\omega} \times \left\{ \left(\frac{d}{dt}\right)_k \vec{r} + \vec{\omega} \times \vec{r} \right\}
$$
\n
$$
\left(\frac{d^2}{dt^2}\right)_l \vec{r} = \left(\frac{d^2}{dt^2}\right)_R \vec{r} + \left(\frac{d}{dt}\right)_R \left(\vec{\omega} \times \vec{r}\right) + \vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r} + \vec{\omega} \times \left(\vec{\omega} \times \vec{r}\right)
$$
\nMultiplying by mass 'm', we shall get quantities that have dimension s's''''b' force'.

Multiplying by mass 'm', we shall get quantities that have dimensions of  $\frac{1}{2}$  of  $\frac{1}{2}$  orce'.

$$
m\left(\frac{d^2}{dt^2}\right)_R \vec{r} = m\left(\frac{d^2}{dt^2}\right)_I \vec{r} - m\left(\frac{d\vec{\omega}}{dt}\right)_R \vec{r} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r} +
$$

$$
- m\vec{\omega} \times (\vec{\omega} \times \vec{r})
$$

$$
\vec{F}_R = \vec{F}_I - \vec{F}_{\omega} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r} - m\vec{\omega} \times \left(\vec{\omega} \times \vec{r}\right)
$$
\n
$$
\begin{array}{|c|c|}\n\hline\n\text{`Leap second' term} & \text{`Coriolis force'} & \text{`Centrifugal force'} \\
\hline\n\text{`Leap second' term} & \text{`Coriolis force'} & \text{`Centrifugal force'} \\
\hline\n\text{`Georiolis term} & \text{`1792} \text{c} \text{d} \cdot \text{``A6} \\
\hline\n\end{array}
$$

We will take a Break… *…… Any questions ?*

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#### **Next L17 : Coriolis Deflection** (47 60  **Foucault Pendulum Cyclonic storm's direction Real Effects of Pseudo-forces! Real Effects of Pseudo-forces**<br> $\vec{F}_R = \vec{F}_I - \vec{F}_{\dot{\omega}} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_\text{s} \vec{r} \text{ } - \text{m}\vec{\omega} \times \left(\vec{\omega} \times \vec{r}\right)$ *d* **Real Effects of Pseudo-force**<br> $\vec{F}_R = \vec{F}_I - \vec{F}_{\dot{\omega}} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_* \vec{r} \ - \ m\vec{\omega} \times \left(\vec{\omega} \times \vec{r}\right)$ Effects of Pseudo-forces! $\vec{\omega}\times \left(\frac{d}{dt}\right)\ \vec{r}\,-\,\mathrm{m}\vec{\omega}\times\left(\vec{\omega}\times\vec{r}\right)$ cts of Pse $\left(\begin{array}{c} d \ \end{array}\right)\ \vec{r}\ -n$ **Real Effects of Pseudo-forces!**<br>= $\vec{F}_I - \vec{F}_{\dot{\omega}} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r}$  - m $\vec{\omega} \times (\vec{\omega} \times \vec{r})$

 $\omega$   $\omega$   $dt$ 

*R*

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#### STiCM Lecture 17: Unit 5

Non-Inertial Frame of Reference cause', when there really isn't one! Real Effects of Pseudo-forces!



#### **Why are Leap Seconds Used?**

The time taken by the earth to do one rotation differs from day to day and from year to year.

The Earth was slower than atomic clocks by 0.16 seconds in 2005;

> by 0.30 seconds in 2006; by 0.31 seconds in 2007; and by 0.32 seconds in 2008.

It was only 0.02 seconds slower in 2001.

http://www.timeanddate.com/time/leapseconds.html ; 14th October, 2009

The atomic clocks can be reset to add an extra second, known as the leap second, to synchronize the atomic clocks with the Earth's observed rotation.

The most accurate and stable time comes from atomic clocks but for navigation and astronomy purposes, *but it is the atomic time that is synchronized with the Earth's rotation.*

#### 'Leap second' term

The International Earth Rotation and Reference System Service (IERS) decides when to introduce a leap second in UTC (**C**oordinated **U**niversal **T**ime).

On one average day, the difference between atomic clocks and Earth's rotation is around 0.002 seconds, or around 1 second every 1.5 years.

IERS announced on July 4, 2008, that a leap second would be added at 23:59:60 (or near midnight) UTC on December 31, 2008. This was the 24th leap second to be added since the first leap second was added in 1972.

http://www.timeanddate.com/time/leapseconds.html ; 14th October, 2009 PCD\_STiCM

## Let us consider 3 'definitions' of the 'vertical'

[1] 'vertical' is defined by the radial line from the center of the earth to a point on the earth's surface



[2] 'vertical' i[s defined by a 'plumb line'](http://z.about.com/d/space/1/0/X/e/gpn-2000-001138.jpg)  suspended at the point under consideration.

[3] 'vertical' is defined by the space curve along which a chalk falls, if you let it go!

# 3 'definitions' of the 'vertical'



$$
m\left(\frac{d^2}{dt^2}\right)_R \vec{r} = m\left(\frac{d^2}{dt^2}\right)_I \vec{r} - m\left(\frac{d\vec{\omega}}{dt}\right)_R \times \vec{r} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r}
$$

$$
- m\vec{\omega} \times (\vec{\omega} \times \vec{r})
$$

#### 7.2921159 × 10−5 radians per second

$$
-m\left(\frac{d\vec{\omega}}{dt}\right)_{R} \vec{r}
$$

$$
m\left(\frac{d^2}{dt^2}\right)_R \vec{r} = m\left(\frac{d^2}{dt^2}\right)_I \vec{r} - m\left(\frac{d\vec{\omega}}{dt}\right)_R \times \vec{r} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r}
$$

$$
- m\vec{\omega} \times \left(\vec{\omega} \times \vec{r}\right)_{\text{http://commons.wikipedia.org/wiki/File:Globespin,git}}
$$

## Earth and moon seen from MARS



First image of Earth & Moon, ever taken from another planet, from Mars, by MGS, on 8 May 2003 at 13:00 GMT.

**MARS GLOBAL SURVEYOR** 

http://mars8.jpl.nasa.gov/mgs/

http://space.about.com/od/pictures/ig/Earth-Pictures-Gallery/Earth-and-Moon-asviewed-from-.htm <u>Paranthe</u>

The moon is attracted toward the earth due to their mutual gravitational attraction.

Will there be an eventual collision?

If so, when? If not, why not?





[1] When travelling in a car/bus, you experience what you call as a 'centrifugal force'. Which physical agency exerts it ? Is there a corresponding force of reaction?

[2] How does a centrifuge work ?

[3] Do the centripetal and the centrifugal forces constitute an 'action - reaction' pair of Newton's 3rd law?



#### **" Foo-Koh"**

The plane of oscillation of the Foucault pendulum is seen to rotate due to the Coriolis effect.

The plane rotates through one full rotation in 24 hours at poles, and in ~33.94 hours at a latitude of 45<sup>0</sup> (Latitude of Paris is  $~190$ .



$$
\overrightarrow{F}_R = \overrightarrow{F}_I - \overrightarrow{F}_{\omega} - 2m\overrightarrow{\omega} \times \left(\frac{d}{dt}\right)_R \overrightarrow{r} - m\overrightarrow{\omega} \times \left(\overrightarrow{\omega} \times \overrightarrow{r}\right)
$$
  
 'Leap second' term 'Coriolis force' 'Centrifugal force'



Gaspard Gustave de **Coriolis** 



1792 - 1843 Jean Bernard Léon Foucault 1819 - 1868

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$$
m\left(\frac{d^2}{dt^2}\right)_R \vec{r} = m\left(\frac{d^2}{dt^2}\right)_I \vec{r}
$$

$$
-m\left(\frac{d\vec{\omega}}{dt}\right)_R \times \vec{r}
$$

$$
-2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r}
$$

$$
\left(\frac{-m\vec{\omega} \times (\vec{\omega} \times \vec{r})}{2m\vec{\omega} \times (\vec{\omega} \times \vec{r})}\right)
$$

We often ignore the 'leap second' and the centrifugal term.





Terrestrial equatorial radius is 6378 km.

Polar radius is 6357ckm PCDS\$TiQM

$$
\vec{F}_{Coriolis} = -2m\vec{\omega} \times \left(\frac{d}{dt}\right)_{R} \vec{r} \quad \vec{\omega} = (\vec{\omega} \cdot \hat{e}_{y})\hat{e}_{y} + (\vec{\omega} \cdot \hat{e}_{z})\hat{e}_{z}
$$
\n
$$
\vec{e}_{z} = \hat{e}_{y}
$$
\n
$$
\vec{\omega}_{z}
$$
\n
$$
\vec{e}_{x} = \hat{e}_{z}
$$
\n
$$
\vec{\omega}_{z}
$$
\n
$$
\vec{e}_{x} = \hat{e}_{z}
$$
\n
$$
\vec{\omega}_{z}
$$
\n
$$
\vec{e}_{z} = \hat{e}_{y} \times \hat{e}_{z}
$$
\n
$$
\vec{\omega}_{z}
$$
\n
$$
\vec{\omega}_{z
$$

$$
\vec{F}_{Coriolis} = -2m\vec{\omega} \times \left(\frac{d}{dt}\right)_{R} \vec{r} \quad \vec{\omega} = (\vec{\omega} \cdot \hat{e}_{y})\hat{e}_{y} + (\vec{\omega} \cdot \hat{e}_{z})\hat{e}_{z}
$$
\n
$$
\hat{e}_{z} = \hat{e}_{v} \text{ (local vertical)}
$$
\n
$$
\hat{e}_{y} = \hat{e}_{\text{North}}
$$
\n
$$
\vec{\omega} \cdot \vec{\omega}
$$
\n
$$
\vec{\omega} \
$$

$$
\vec{F}_{Coriolis} = -2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r}
$$

$$
\vec{\omega} = (\vec{\omega} \cdot \hat{e}_y) \hat{e}_y + (\vec{\omega} \cdot \hat{e}_z) \hat{e}_z
$$

Coriolis deflection of an object in 'free' 'fall' at a point on earth's surface:

$$
\left(\frac{d}{dt}\right)_R \vec{r} = \vec{v}_R = v_R(-\hat{e}_z)
$$

nt on earth's surface:  
\n
$$
\vec{F}_{Coriolis} = -2m \left[ \left( \vec{\omega} \bullet \hat{e}_y \right) \hat{e}_y + \left( \vec{\omega} \bullet \hat{e}_z \right) \hat{e}_z \right] \times \mathbf{v}_R (-\hat{e}_z)
$$
\n
$$
= 2m \left( \vec{\omega} \bullet \hat{e}_y \right) \hat{e}_x = 2m \omega \cos \lambda \hat{e}_{East}
$$

 $0 \leq \lambda \leq \frac{\pi}{2}$  ( 0 (S hemisphere) 2 ) *N hemisphere* 2 *S hemisphere*  $\pi$  $\boldsymbol{\lambda}$  $\pi$  $-\frac{\pi}{2} \leq \lambda \leq 0$  (  $\leq \lambda \leq \frac{\pi}{2}$  (*N* hemisphere) Coriolis deflection: toward East in both the Northern & *here*) the Southern Hemispheres<br> $\cos \lambda$ <sup>PCD</sup>-<sup>STICM</sup> oth N- and S-hemispheres PCD\_STiCM



#### http://www.youtube.com/watch?v=ZvLYrZ3vgio&NR=1



This Nanyuki tourist attraction is *not* due to Coriolis effect









# What will be the effect on rockets?



On ICBMs? 'Missile Woman of India' Dr. Tessy Thomas



PCD\_STiQMTh May, 2010 AGNI II
http://daphne.palomar.edu/pdeen/Animations/34\_Coriolis.swf



We will take a Break… *…… Any questions ?*

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# **Next L18 : Coriolis Deflection Foucault Pendulum Real Effects of Pseudo-forces!**

$$
\vec{F}_R = \vec{F}_I - \vec{F}_{\dot{\omega}} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r} - m\vec{\omega} \times \left(\vec{\omega} \times \vec{r}\right)
$$

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STICM Lecture 18: Unit 5

'EFFECT', when there isn't a cause! 'CAUSE', when there really isn't one! Real Effects of Pseudo-forces! Foucault Pendulum

$$
\vec{F}_R = \vec{F}_I - \vec{F}_{\dot{\omega}} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r} - m\vec{\omega} \times (\vec{\omega} \times \vec{r})
$$



*Physical Oceanography of the Baltic Sea Matti Leppäranta, Kai Myrberg Springer-Praxis (2009)*

**Circulating** current in the Baltic sea (1969?)



### Coriolis effect – only for Physicists?

08

0

Electronics engineers? Computer science engineers? Aerospace engineers? **Ocean Engineering?** 

Any navigation system on earth..... GPS in cellphones!

器の目 用る

Use a Cartesian coordinate system with reference to a point on the earth's surface.

Choose  $\left\{\hat{e}_x,\hat{e}_y,\hat{e}_z\right\}$  such that  $\hat{e}_z$  is along the local 'up/vertical' direction (which is not along  $\hat{e}^{\phantom{\dag}}_{r}$  due to the centrifugal term).

Choose  $\hat{e}_y$  such that it is orthogonal to  $\hat{e}_z$ , and points toward the North-pole seen from the point on the earth's surface under consideration. ,  $\hat{e}_z$  *is very neal*<br>uch that it is orthogo<br>orth-pole seen from the consideration.<br>the consideration.<br>, choose  $\hat{e}_x = \hat{e}_y \times d$ centrifugal term).<br>*However,*  $\hat{e}_z$  *is very nearly the same as*  $\hat{e}_r$ 

Finally, choose  $\hat{e}_x = \hat{e}_y \times \hat{e}_z$ , which will give us the direction of the local 'East' at that point.

$$
\vec{F}_{Coriolis} = -2m\vec{\omega} \times \left(\frac{d}{dt}\right)_{R} \vec{r} \quad \vec{\omega} = (\vec{\omega} \cdot \hat{e}_{y})\hat{e}_{y} + (\vec{\omega} \cdot \hat{e}_{z})\hat{e}_{z}
$$
\n
$$
\lambda = \angle (\vec{\omega}, \hat{e}_{North}) = \angle (\vec{\omega}, \hat{e}_{y}) \qquad \hat{e}_{z} = \hat{e}_{y}
$$
\n
$$
\lambda = \frac{\pi}{2} \text{ at N Pole}
$$
\n
$$
\lambda = -\frac{\pi}{2} \text{ at Spole}
$$
\n
$$
\lambda = -\frac{\pi}{2} \text{ at Spole}
$$
\n
$$
\hat{e}_{x} = \hat{e}_{\text{fast}} = \hat{e}_{y} \times \hat{e}_{z}
$$
\n
$$
\hat{e}_{z} = \hat{e}_{y} \text{ (local vertical')}
$$
\n
$$
\hat{e}_{y} = \hat{e}_{North}
$$
\n
$$
\hat{e}_{x} = \hat{e}_{East} = \hat{e}_{y} \times \hat{e}_{z}
$$

Caution! We use a mix of three coordinate systems!

- 1. A cartesian coordinate system as defined in the previous slide.
- 2. A spherical polar coordinate system whose 'polar' angle is defined with respect to the axis of the earth's rotation rather than with respect to the cartesian z-axis which is oriented along the local 'vertical'.
- 3. Cylindrical polar coordinate system whose radial unit vector is along the radial outward direction with reference to the earth's <u>axis</u> of rotation.

Just follow the 6 steps indicated in the next slide, exactly in the order given!



rotation! The 'vertical' is *not* along the radial line, nor ˆ *F F e cf cf* i<br>L along the axis of earth's rotation!

$$
\text{PCD\_STIC}(\vec{w}) = (\vec{\omega} \cdot \hat{e}_y) \hat{e}_y + (\vec{\omega} \cdot \hat{e}_z) \hat{e}_z
$$

$$
\vec{F}_R = \vec{F}_I - \vec{F}_{\omega} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r} - m\vec{\omega} \times (\vec{\omega} \times \vec{r})
$$
\n
$$
\vec{\omega} = (\vec{\omega} \cdot \hat{e}_y) \hat{e}_y + (\vec{\omega} \cdot \hat{e}_z) \hat{e}_z
$$
\n
$$
\left(\frac{d}{dt}\right)_R \vec{r} = \left[v_x \hat{e}_x + v_y \hat{e}_y + v_z \hat{e}_z\right]
$$
\nVelocity of the object  
\nROTATING FRAME

$$
\vec{\omega} = (\vec{\omega} \cdot \hat{e}_y) \hat{e}_y + (\vec{\omega} \cdot \hat{e}_z) \hat{e}_z
$$

$$
\left(\frac{d}{dt}\right)_R = \left[v_x \hat{e}_x + v_y \hat{e}_y + v_z \hat{e}_z\right]
$$
\n
$$
\lambda : latitude
$$

Velocity of the object in ROTATING FRAME



$$
\lambda: latitude
$$

$$
\overrightarrow{F}_{Coriolis} = -2m \left[ \left( \vec{\omega} \bullet \hat{e}_y \right) \hat{e}_y + \left( \vec{\omega} \bullet \hat{e}_z \right) \hat{e}_z \right] \times \left[ v_x \hat{e}_x + v_y \hat{e}_y + v_z \hat{e}_z \right]
$$
\n
$$
FCD\_STICM
$$

PCD\_STiCM

$$
\vec{F}_{Coriolis} = -2m\omega \left[ \cos \lambda \hat{e}_y + \sin \lambda \hat{e}_z \right] \times \left[ \mathbf{v}_x \hat{e}_x + \mathbf{v}_y \hat{e}_y + \mathbf{v}_z \hat{e}_z \right]
$$

$$
m\vec{a}_{Coriolis} = -2m\omega \left[ \left( \cos \lambda v_z - \sin \lambda v_y \right) \hat{e}_x + \sin \lambda v_x \hat{e}_y + (-\cos \lambda v_x) \hat{e}_z \right]
$$



### Pendulum at the N pole

## The observer is at the N pole. The pole of the pendulum  $\mathbb{R}^n$

PCD\_STiCM

look to an observer at the control of the control of

N pole, but standing

disconnected with the state of the

How will the plane of oscillation of the pendulum look to an observer at the N pole?

## The observer is at the S pole. The pole of the pendulum of the pole of the pendulum  $\mathcal{P}$ look to an observer at the second state of the second state of the second state of the second state of the second

How will the plane of oscillation of the pendulum look to an observer at the S pole?

### Pendulum at the S pole PCD\_STiCM

N pole, but standing

disconnected with the state of the

## What if the pendulum is at an intermediate latitude?



How will the plane of oscillation of the pendulum look to an observer at that latitude?

Northern edge of the floor of the room is closer to the earth's axis than the Southern edge.

**Northern** edge of the floor moves eastward slower than the Southern edge.

W.B.Somerville Q.Jl. Royal Astronomical Soc. (1972) 13 40-62

**POD\_STICM** 

http://www.youtube.com/watch?v=nB2SXLYwKkM

video compilation of a Foucault pendulum in action at the Houston Museum of Natural Science.

$$
\vec{\omega} = (\vec{\omega} \cdot \hat{e}_y) \hat{e}_y + (\vec{\omega} \cdot \hat{e}_z) \hat{e}_z
$$

$$
\begin{aligned}\n\text{UCault Pendulum} \quad \vec{\omega} &= (\vec{\omega} \cdot \hat{e}_y) \hat{e}_y + (\vec{\omega} \cdot \hat{e}_z) \hat{e}_z \\
\vec{F}_R &= \vec{F}_I - \vec{F}_{\hat{\omega}} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r} - m\vec{\omega} \times (\vec{\omega} \times \vec{r}) \\
m\vec{r}_R &= m\vec{g} + \vec{S} - \vec{P}_{\hat{\phi}} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r} - m\vec{\omega} \times (\vec{\omega} \times \vec{r}) \\
\vec{S} &= \vec{S} \vec{u} \\
\vec{S} &= \vec{S} \vec{u} \\
\vec{S} &= \vec{S} \left[ \hat{e}_x (\hat{e}_x \cdot \vec{\omega}) + \hat{e}_y (\hat{e}_y \cdot \vec{\omega}) + \hat{e}_z (\hat{e}_z \cdot \vec{\omega}) \right] \\
m\vec{g} &= S \left[ \hat{e}_x (\hat{e}_x \cdot \vec{\omega}) + \hat{e}_y (\hat{e}_y \cdot \vec{\omega}) + \hat{e}_z (\hat{e}_z \cdot \vec{\omega}) \right] \\
\vec{S} &= S \left[ \hat{e}_x \cos \alpha + \hat{e}_y \cos \beta + \hat{e}_z \cos \gamma \right]\n\end{aligned}
$$

For example, the following equations is:

\n
$$
\vec{ma}_{\text{Coriolis}} = -2m\omega \left[ \left( \cos \lambda \vec{x}_z - \sin \lambda v_y \right) \hat{e}_x + \sin \lambda v_x \hat{e}_y + \left( -\cos \lambda v_x \right) \hat{e}_z \right]
$$
\n
$$
\vec{S} = S \left[ \hat{e}_x \cos \alpha + \hat{e}_y \cos \beta + \hat{e}_z \cos \gamma \right] \quad \boxed{\text{neglect } \dot{z}}
$$
\n
$$
x, y \text{ motion:}
$$
\n
$$
m\ddot{x} = S \cos \alpha - 2m\omega \left( \cos \lambda \dot{z} - \sin \lambda \dot{y} \right)
$$
\n
$$
\vec{S} = \vec{S} \left[ \hat{e}_x \cos \beta - 2m\omega \left( \cos \lambda \dot{z} - \sin \lambda \dot{y} \right) \right]
$$
\n
$$
m\ddot{y} = S \cos \beta - 2m\omega \sin \lambda \dot{x}
$$
\n
$$
m\ddot{y} = mg \cos \alpha + 2m\omega \sin \lambda \dot{y}
$$
\n
$$
m\ddot{y} = mg \cos \beta - 2m\omega \sin \lambda \dot{x}
$$

 $\cdot$ 

$$
m\ddot{\vec{r}}_R = m\vec{g} + \vec{S} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r}
$$

ult Pendulum 
$$
m\ddot{\vec{r}}_R = m\vec{g} + \vec{S} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r}
$$
  
\n
$$
\frac{\vec{S} = S\hat{u}}{m\ddot{x}} = mg \cos \alpha + 2\eta \omega \sin \lambda \dot{y}
$$
\n
$$
\vec{r} = \eta \cos \beta - 2\eta \omega \sin \lambda \dot{x}
$$
\n
$$
\vec{x} = g \cos \alpha + 2\omega \sin \lambda \dot{y}
$$
\n
$$
\vec{S} = \frac{\vec{y} = g \cos \beta - 2\omega \sin \lambda \dot{x}}{\vec{y} = g \cos \alpha + 2\Omega \dot{y}}
$$
\n
$$
\vec{r} = g \cos \alpha + 2\Omega \dot{y}
$$
\n
$$
\vec{y} = g \cos \beta - 2\Omega \dot{x}
$$

$$
\ddot{x} = g \cos \alpha + 2\Omega \dot{y}
$$
\nCoupled differential

\n
$$
\ddot{y} = g \cos \beta - 2\Omega \dot{x}
$$
\nequations

such that

Solve by transforming to new coordinates x',y'  
such that 
$$
\begin{bmatrix} x = x' \cos(\Omega t) + y' \sin(\Omega t) \\ y = -x' \sin(\Omega t) + y' \cos(\Omega t) \end{bmatrix}
$$

$$
y = -x \sin(32t) + y \cos(32t)
$$
  
\n
$$
\dot{x} = -\Omega \left[ x \sin(\Omega t) - y \cos(\Omega t) \right]
$$
  
\n
$$
\dot{y} = -\Omega \left[ x \cos(\Omega t) + y \sin(\Omega t) \right]
$$
  
\n
$$
\ddot{x} = g \cos \alpha - 2\Omega^2 \left[ x \cos(\Omega t) + y \sin(\Omega t) \right]
$$
  
\n
$$
\ddot{y} = g \cos \beta + 2\Omega^2 \left[ x \sin(\Omega t) - y \cos(\Omega t) \right]
$$

$$
\ddot{x} = g \cos \alpha - 2\Omega^2 \left[ x' \cos(\Omega t) + y' \sin(\Omega t) \right]
$$
  

$$
\ddot{y} = g \cos \beta + 2\Omega^2 \left[ x' \sin(\Omega t) - y' \cos(\Omega t) \right]
$$

 $\vec{x}\cos(\Omega t) = g\cos\alpha\cos(\Omega t) - 2\Omega^2 \left[ x'\cos^2(\Omega t) + y'\sin(\Omega t)\cos(\Omega t) \right]$  $\ddot{y}\sin(\Omega t) = g\cos\beta\sin(\Omega t) + 2\Omega^2 \left[x'\sin^2(\Omega t) - y'\cos(\Omega t)\sin(\Omega t)\right]$ 

 $\vec{x}\cos(\Omega t) + \ddot{y}\sin(\Omega t) = g\cos\alpha\cos(\Omega t) - 2\Omega^2 \left[x'\cos^2(\Omega t) + y'\sin(\Omega t)\cos(\Omega t)\right]$ +g cos  $\beta$  sin  $(\Omega t)$  +  $2\Omega^2 \left[ x \sin^2 (\Omega t) - y \cos (\Omega t) \sin (\Omega t) \right]$ 

$$
\begin{aligned} \n\ddot{x}\cos(\Omega t) + \ddot{y}\sin(\Omega t) &= g\cos\alpha\cos(\Omega t) - 2\Omega^2 \left[ x'\cos^2(\Omega t) + y'\sin(\Omega t)\cos(\Omega t) \right] \\ \n&\quad + g\cos\beta\sin(\Omega t) + 2\Omega^2 \left[ x'\sin^2(\Omega t) - y'\cos(\Omega t)\sin(\Omega t) \right] \n\end{aligned}
$$

Dropping terms in  $\Omega^2$  $(\ddot{x} - g \cos \alpha)\cos(\Omega t) + (\ddot{y} - g \cos \beta)\sin(\Omega t) = 0$  $(\ddot{x} - g \cos \alpha) = 0$  $(\ddot{y} - g \cos \beta) = 0$ 

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$$
\Omega = \omega \sin \lambda
$$

$$
m\ddot{\vec{r}}_R = m\vec{g} + \vec{S} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r}
$$

Direction cosines of  $\vec{S}$  are:

$$
(\ddot{x} - g \cos \alpha) = 0
$$
  

$$
(\ddot{y} - g \cos \beta) = 0
$$

 $\ddot{x} + g$ 

 $\ddot{y} + g$ 

$$
\Omega = \omega \sin \lambda
$$
\n $m\vec{r}_R = m\vec{g} + S - 2m\omega \times$ \n
$$
\begin{aligned}\n&\text{Direction cosines of} \\
&(\vec{x} - g \cos \alpha) = 0 \\
&(\vec{y} - g \cos \beta) = 0\n\end{aligned}
$$
\n
$$
\begin{aligned}\n&\text{Direction cosines of} \\
&\cos \alpha = -\frac{x}{l} \\
&\cos \beta = -\frac{y}{l}\n\end{aligned}
$$
\n
$$
\begin{aligned}\n&\text{cos } \gamma = \frac{z - l}{l} \\
&(\vec{y} + g \frac{y}{l}) = 0\n\end{aligned}
$$
\n
$$
\begin{aligned}\n&\text{Two-dimensional linear form} \\
&\text{In a linear combination of the form } \alpha = \frac{z - l}{l}\n\end{aligned}
$$

Two-dimensional linear

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0

*x*

*y*

 $\left(\ddot{y} + g\frac{y}{l}\right) = 0$ 

 $\left(\ddot{x} + g\frac{x}{l}\right) = 0$ 

0

$$
m\ddot{\vec{r}}_R = m\vec{g} + \vec{S} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r}
$$

*S* endulum  $m\ddot{\vec{r}}_R = m\vec{g} + \vec{S} - 2m\vec{\omega}\times \left(\frac{d}{dt}\right)_R \vec{r}$ <br>
erms in  $\Omega^2$ <br>  $\left(\cos(\Omega t) + (\ddot{y} - g \cos \beta)\sin(\Omega t) = 0\right)$ <br>  $\left(\ddot{x} + g\frac{x}{l}\right) = 0; \quad \left(\ddot{y} + g\frac{y}{l}\right) = 0$ <br>
in the earth's rotating frame, the path is that<br>
of an e In the earth's rotating frame, the path is that of an ellipse.  $\begin{pmatrix} y & 8 \cos \theta & y \sin \theta & z \end{pmatrix} = 0; \quad \begin{pmatrix} \ddot{y} + g & \frac{y}{z} \end{pmatrix} = 0$  $\ddot{x} + g\frac{x}{l} = 0; \quad (\ddot{y} + g)$  $\left(\frac{x}{l}\right) = 0; \quad \left(\ddot{y} + g\frac{y}{l}\right)$  $(\ddot{x}-g\cos\alpha)\cos(\Omega t)+(\ddot{y}-g\cos\beta)\sin(\Omega t)=0$ <br>  $(\ddot{x}+g\frac{x}{l})=0; \quad (\ddot{y}+g\frac{y}{l})=0$ Dropping terms in  $\Omega^2$ Propping terms in  $\Omega^2$ <br>  $\ddot{x} - g(\cos \alpha) \cos(\Omega t) + (\ddot{y} - g \cos \beta) \sin(\Omega t) = 0$  $\boldsymbol{\Omega}$ opping terms in  $\Omega^2$ <br>-  $g$  cos  $\alpha$  cos  $(\Omega t) + (\ddot{y} - g \cos \beta) \sin (\Omega t) = 0$ 

$$
m\ddot{\vec{r}}_R = m\vec{g} + \vec{S} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r}
$$

*S* endulum  $m\ddot{\vec{r}}_R = m\vec{g} + \vec{S} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r}$ <br>
rms in  $\Omega^2$ <br>  $(\cos(\Omega t) + (\ddot{y} - g \cos \beta) \sin(\Omega t) = 0$ <br>  $+ g \frac{x}{l} = 0; \quad (\ddot{y} + g \frac{y}{l}) = 0$ <br>
The ellipse would precess at an angular<br>
speed  $\Omega = \omega \sin \lambda$ <br>
A number of a The ellipse would precess at an angular speed  $\Omega = \omega \sin \lambda$ A number of approximations made! A number of approximations made!<br> $m\vec{g}$  Detailed analysis is rather involved!  $(\Omega t) + (\ddot{y} - g \cos \beta) \sin \left(\frac{x}{y}\right) = 0; \quad (\ddot{y} + g \frac{y}{x}) = 0$  $\ddot{x} + g \frac{x}{l} = 0; \quad (\ddot{y} + g)$  $(\ddot{x} - g\cos\alpha)\cos(\Omega t) + (\ddot{y} - g\cos\beta)\sin(\Omega t)$ <br>  $(\ddot{x} + g\frac{x}{l}) = 0; \quad (\ddot{y} + g\frac{y}{l}) = 0$ Dropping terms in  $\Omega^2$ Propping terms in  $\Omega^2$ <br>  $\ddot{x} - g(\cos \alpha) \cos(\Omega t) + (\ddot{y} - g \cos \beta) \sin(\Omega t) = 0$ The "plane" would in fact be a "curved surface".

$$
m\ddot{\vec{r}}_R = m\vec{g} + \vec{S} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r}
$$

The ellipse would precess at an angular speed  $\Omega = \omega \sin \lambda$ *S* endulum  $m\ddot{\vec{r}}_R = m\vec{g} + \vec{S} - 2m\vec{\omega} \times \left(\frac{d}{dt}\right)_R \vec{r}$ <br>
ellipse would precess at an angular<br>
ed  $\Omega = \omega \sin \lambda$ <br>  $\lambda : latitude$ <br>
Time Period for the rotation of the<br>
"plane" of oscillation of the Foucault<br>
Pendulum<br>  $T = \frac{1}{f$ : *latitude*  $\frac{1}{2} = \frac{2\pi}{4} = \frac{2}{1}$ sin  $\frac{2\pi}{2\pi} = \frac{24 \text{ hours}}{24 \text{ hours}}$  $\frac{2\pi}{2\pi v \sin \lambda} = \frac{24 \text{ hours}}{\sin \lambda}$  $\frac{\pi}{4} - \frac{2\pi}{4}$  $\frac{2\pi}{\omega \sin \lambda}$  $=\frac{1}{f}=\frac{2\pi}{\Omega}=\frac{2\pi}{\Omega}$  $\Omega$  $=\frac{2\pi}{2\pi i \sin^2 2}=\frac{241}{24}$ *T f* Time Period for the rotation of the "plane" of oscillation of the Foucault Pendulum PCD\_STICM $\pi v \sin \lambda$  Sin  $\lambda$ 



Famous Foucault Pendulum:

- at the Pantheon in Paris, France.

http://www.animations.physics.unsw.edu.a u/jw/foucault\_pendulum.html

#### Wire going up to ceiling

http://www.youtube.com/watch?v=vVg5P6f rHzY&feature=related

Weight on long wire

Disk to show direction PGD<sub>S</sub>



## Philosophical questions: What is 'force' ? Mass/ Inertial frame? Gravity?





Ernst Mach (1838–1916)



1879 – 1955

Newton: Gravity is the result of an attractive interaction between objects having mass. Curvature of Space-Time, Geometry / Dynamics of Matter  $\oint$  General Theory of Relativity

From a portrait by Enoch Seeman in 1726 Sir Isaac Newton 1643 - 1727 Albert Einstein

We will take a Break… *…… Any questions ?*

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• In the next Unit, we shall consider Lorentz Transformations and Einstein's Special Theory of Relativity.

• *c: finite!* 



## **Next, Unit 6: Special Theory of Relativity**

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